

K⁰

$$I(J^P) = \frac{1}{2}(0^-)$$

K⁰ MASS

<u>VALUE (MeV)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
497.648±0.022 OUR FIT				
497.648±0.022 OUR AVERAGE				
497.625±0.001±0.031	655k	LAI	02 NA48	K_L^0 beam
497.661±0.033	3713	BARKOV	87B CMD	$e^+ e^- \rightarrow K_L^0 K_S^0$
497.742±0.085	780	BARKOV	85B CMD	$e^+ e^- \rightarrow K_L^0 K_S^0$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
497.44 ±0.50		FITCH	67 OSPK	
498.9 ±0.5	4500	BALTAY	66 HBC	K^0 from $\bar{p}p$
497.44 ±0.33	2223	KIM	65B HBC	K^0 from $\bar{p}p$
498.1 ±0.4		CHRISTENS...	64 OSPK	

 $m_{K^0} - m_{K^\pm}$

<u>VALUE (MeV)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>CHG</u>	<u>COMMENT</u>
3.972±0.027 OUR FIT		Error includes scale factor of 1.2.			
• • • We do not use the following data for averages, fits, limits, etc. • • •					
3.95 ±0.21	417	HILL	68B DBC	+	$K^+ d \rightarrow K^0 pp$
3.90 ±0.25	9	BURNSTEIN	65 HBC	-	
3.71 ±0.35	7	KIM	65B HBC	-	$K^- p \rightarrow n\bar{K}^0$
5.4 ±1.1		CRAWFORD	59 HBC	+	
3.9 ±0.6		ROSENFIELD	59 HBC	-	

K⁰ MEAN SQUARE CHARGE RADIUS

<u>VALUE (fm²)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
-0.077±0.010 OUR AVERAGE				
• • • We do not use the following data for averages, fits, limits, etc. • • •				
-0.077±0.007±0.011	5037	ABOUZAID	06 KTEV	$K_L^0 \rightarrow \pi^+ \pi^- e^+ e^-$
-0.090±0.021		LAI	03C NA48	$K_L^0 \rightarrow \pi^+ \pi^- e^+ e^-$
-0.054±0.026		MOLZON	78	K_S^0 regen. by electrons
-0.087±0.046		BLATNIK	79	VMD + dispersion relations
-0.050±0.130		FOETH	69B	K_S^0 regen. by electrons

T-VIOLATION PARAMETER IN K⁰- \bar{K}^0 MIXING

The asymmetry $A_T = \frac{\Gamma(\bar{K}^0 \rightarrow K^0) - \Gamma(K^0 \rightarrow \bar{K}^0)}{\Gamma(\bar{K}^0 \rightarrow K^0) + \Gamma(K^0 \rightarrow \bar{K}^0)}$ must vanish if T invariance holds.

ASYMMETRY A_T IN K⁰- \bar{K}^0 MIXING

<u>VALUE (units 10⁻³)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
6.6±1.3±1.0	640k	¹ ANGELOPO...	98E CPLR

¹ ANGELOPOULOS 98E measures the asymmetry $A_T = [\Gamma(\bar{K}_{t=0}^0 \rightarrow e^+ \pi^- \nu_{t=\tau}) - \Gamma(K_{t=0}^0 \rightarrow e^- \pi^+ \bar{\nu}_{t=\tau})]/[\Gamma(\bar{K}_{t=0}^0 \rightarrow e^+ \pi^- \nu_{t=\tau}) + \Gamma(K_{t=0}^0 \rightarrow e^- \pi^+ \bar{\nu}_{t=\tau})]$ as a function of the neutral-kaon eigentime τ . The initial strangeness of the neutral kaon is tagged by the charge of the accompanying charged kaon in the reactions $p\bar{p} \rightarrow K^-\pi^+K^0$ and $p\bar{p} \rightarrow K^+\pi^-\bar{K}^0$. The strangeness at the time of the decay is tagged by the lepton charge. The reported result is the average value of A_T over the interval $1\tau_S < \tau < 20\tau_S$. From this value of A_T ANGELOPOULOS 01B, assuming *CPT* invariance in the $e\pi\nu$ decay amplitude, determine the *T*-violating as $\Delta S = \Delta S$ conserving parameter (for its definition, see Review below) $4\text{Re}(\epsilon) = (6.2 \pm 1.4 \pm 1.0) \times 10^{-3}$.

CPT INVARIANCE TESTS IN NEUTRAL KAON DECAY

Revised June 2006 by P. Bloch (CERN).

The time evolution of a neutral kaon state state is described by

$$\frac{d}{dt}\Psi = -i\Lambda\Psi, \quad \Lambda \equiv M - \frac{i}{2}\Gamma \quad (1)$$

where M and Γ are Hermitian 2×2 matrices known as the mass and decay matrices. The corresponding eigenvalues are $\lambda_{L,S} = m_{L,S} - \frac{i}{2}\gamma_{L,S}$. *CPT* invariance requires the diagonal elements of Λ to be equal. The *CPT*-violation complex parameter δ is defined as

$$\begin{aligned} \delta &= \frac{\Lambda_{\bar{K}^0\bar{K}^0} - \Lambda_{K^0K^0}}{2(\lambda_L - \lambda_S)} \\ &= \delta_{\parallel} \exp(i\phi_{SW}) + \delta_{\perp} \exp(i(\phi_{SW} + \frac{\pi}{2})) \end{aligned} \quad (2)$$

where we have introduced the projections δ_{\parallel} and δ_{\perp} respectively parallel and perpendicular to the superweak direction $\phi_{SW} = \tan^{-1}(2\Delta m/\Delta\gamma)$, where $\Delta m = m_L - m_S$ and $\Delta\gamma = \gamma_S - \gamma_L$, the positive mass and width differences between K_L and K_S . These projections are linked to the mass and width difference between K^0 and \bar{K}^0 :

$$\delta_{\parallel} = \frac{1}{4} \frac{\gamma_{K^0} - \gamma_{\bar{K}^0}}{\sqrt{\Delta m^2 + \left(\frac{\Delta\gamma}{2}\right)^2}}, \quad \delta_{\perp} = \frac{1}{2} \frac{m_{K^0} - m_{\bar{K}^0}}{\sqrt{\Delta m^2 + \left(\frac{\Delta\gamma}{2}\right)^2}}. \quad (3)$$

$\text{Re}(\delta)$ can be directly measured by studying the time evolution of the strangeness content of initially pure K^0 and \bar{K}^0 states, for example through the asymmetry

$$A_{CPT} = \frac{P[\bar{K}^0 \rightarrow \bar{K}^0(t)] - P[K^0 \rightarrow K^0(t)]}{P[\bar{K}^0 \rightarrow \bar{K}^0(t)] + P[K^0 \rightarrow K^0(t)]} = 4\text{Re}(\delta) \quad (4)$$

where $P[a \rightarrow b(t)]$ is the probability that the pure initial state a is seen as state b at proper time t . This method has been used by tagging the initial strangeness with strong interactions and the final strangeness with the semileptonic decay (a more appropriate combination of semileptonic rates allows to be independent of any direct CPT violation in the decay itself) and yields today's best value of $\text{Re}(\delta)$, compatible with zero with an error of $\sim 3 \times 10^{-4}$.

As an alternative it has been proposed to compare the semileptonic charge asymmetries for K_L and K_S

$$A_{L,S} = \frac{R(K_{L,S} \rightarrow \pi^- \ell^+ \nu) - R(K_{L,S} \rightarrow \pi^+ \ell^- \bar{\nu})}{R(K_{L,S} \rightarrow \pi^- \ell^+ \nu) + R(K_{L,S} \rightarrow \pi^+ \ell^- \bar{\nu})},$$

$$A_S - A_L = 4\text{Re}(\delta). \quad (5)$$

A_L has been accurately measured. A_S has been recently measured with tagged K_S at ϕ factories, however not yet with the required accuracy. Note however that Eq. (5) assumes CPT invariance in the $\Delta S = -\Delta Q$ semileptonic decay amplitude.

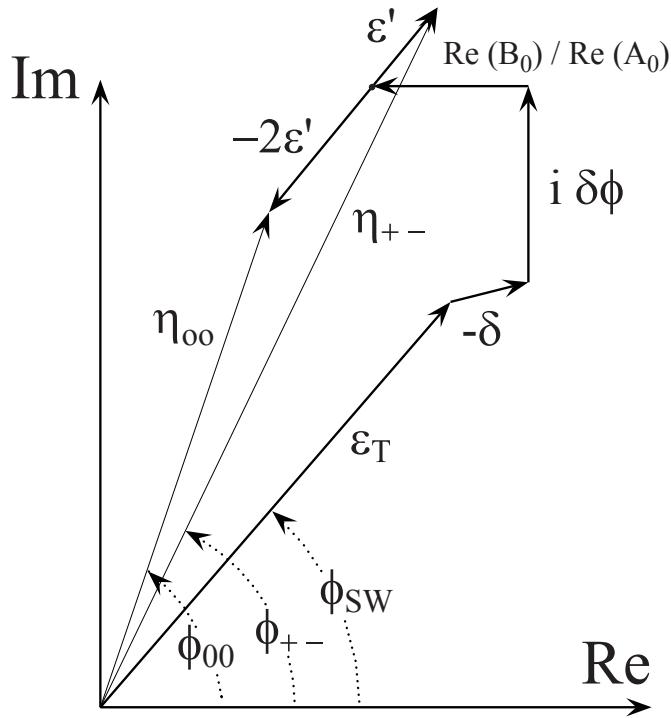


Figure 1: CP - and CPT -violation parameters in 2π decay.

δ_\perp can be obtained from the measurement of the $\pi\pi$ decays CP -violation parameters η_{+-} and η_{00} . Figure 1 shows the various contributions to $\eta_{\pi\pi}$ [1]. The T -violation parameter ϵ_T

$$\epsilon_T = i \frac{|\Lambda_{K^0\bar{K}^0}|^2 - |\Lambda_{\bar{K}^0K^0}|^2}{\Delta\gamma(\lambda_L - \lambda_S)} \quad (6)$$

has been defined in such a way that it is exactly aligned along the superweak direction [‡]. A_I (resp. B_I) is the CPT -conserving (resp. violating) decay amplitude for the $\pi\pi$ Isospin I state, ε' is the direct CP/CPT -violation parameter [$\varepsilon' = 1/3(\eta_{+-} - \eta_{00})$] and $\delta\phi = \frac{1}{2} [\varphi_\Gamma - \arg(A_0^*\bar{A}_0)]$ is the phase difference between

the $I = 0$ component of the decay amplitude and the matrix element $\Gamma_{K^0\overline{K}^0}$. From Fig. 1 one obtains

$$\begin{aligned}\delta_{\perp} = & |\eta_{+-}|(\phi_{SW} - \frac{2}{3}\phi_{+-} - \frac{1}{3}\phi_{00}) \\ & - \frac{\text{Re}(B_0)}{\text{Re}(A_0)} \sin(\phi_{SW}) + \delta\phi \cos(\phi_{SW}) .\end{aligned}\quad (7)$$

The present accuracy on the term $|\eta_{+-}|(\phi_{SW} - \frac{2}{3}\phi_{+-} - \frac{1}{3}\phi_{00})$ is 2.6×10^{-5} . $\delta\phi$ gets contributions from CP violation in semileptonic and 3π decays [2,3] and can only be neglected at the present time if one assumes that η_{000} is not significantly larger than η_{+-0} . Furthermore, B_0 is not directly measured, so additional assumptions (for example, CPT conservation in the decay which implies $B_0 = 0$) or a combination with other measurements are necessary to obtain δ_{\perp} .

If one assumes unitarity, one can measure $\text{Im}(\delta)$ using the Bell-Steinberger relation which relates K_S and K_L decay amplitudes into all final states f :

$$\text{Re}(\epsilon_T) - i\text{Im}(\delta) = \frac{1}{2(i\Delta m + \frac{1}{2}(\gamma_L + \gamma_S))} \times \sum A_{fL} A_{fS}^* . \quad (8)$$

Since the $\pi\pi$ amplitudes dominate, the result relies also strongly on the $\phi_{\pi\pi}$ phase measurements. The advantage is that B_0 does not enter. Using all available data, one obtains a value of $\text{Im}(\delta)$ compatible with zero with a precision of 2×10^{-5} . The precision here is limited by the measurement of η_{+-} .

The results on $\text{Re}(\delta)$ and $\text{Im}(\delta)$ can be combined to obtain δ_{\parallel} and δ_{\perp} and therefore the $K^0\overline{K}^0$ mass and width difference shown in Fig. 2. The current accuracy is a few 10^{-18} GeV for both.

If one assumes that CPT is conserved in the decays ($\gamma_{K^0} = \gamma_{\overline{K}^0}$, $\delta_{\parallel} = 0$, $B_I = 0$), the phase of δ is known, and the δ_{\perp} and

Bell-Steinberger methods are identical. One in this case obtains a limit for $|m_{K^0} - m_{\bar{K}^0}|$ of 4.7×10^{-19} GeV (90%CL).

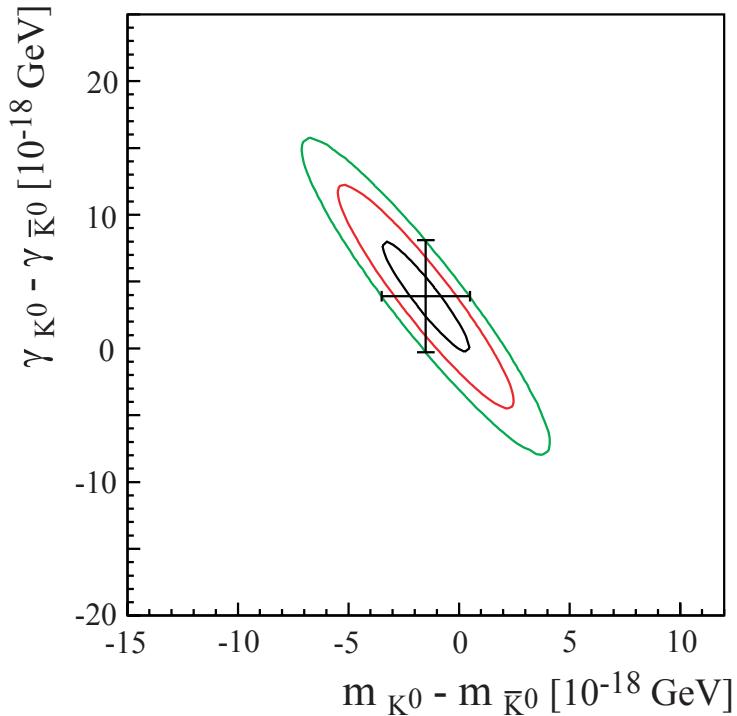


Figure 2: $K^0 - \bar{K}^0$ mass vs width difference.

Footnotes and References

- [‡] Many authors have a different definition of the T -violation parameter, $\epsilon = (\Lambda_{\bar{K}^0 K^0} - \Lambda_{K^0 \bar{K}^0})/(2(\lambda_L - \lambda_S))$. ϵ is not exactly aligned with the superweak direction. The two definitions can be related through $\epsilon = \epsilon_T + i\delta\phi$.
1. See for instance, C.D. Buchanan *et al.*, Phys. Rev. **D45**, 4088 (1992). See also the Second Daphne Handbook, Ed. L.Maiani *et al.*, INFN Frascati (1995).
 2. V.V. Barmin *et al.*, Nucl. Phys. **B247**, 293 (1984).
 3. L. Lavoura, Mod. Phys. Lett. **A7**, 1367 (1992).

CP-VIOLATION PARAMETERS

Re(ϵ)

VALUE (units 10^{-3})	DOCUMENT ID	TECN
1.664±0.010	2 LAI	05A NA48

² LAI 05A values are obtained through unitarity (Bell-Steinberger relations), improving determination of η_{000} and combining other data from PDG and APOSTOLAKIS 99B.

CPT-VIOLATION PARAMETERS

In K^0 - \bar{K}^0 mixing, if *CPT*-violating interactions include a *T* conserving part then

$$|K_S\rangle = [|K_1\rangle + (\epsilon + \delta)|K_2\rangle]/\sqrt{1+|\epsilon+\delta|^2}$$

$$|K_L\rangle = [|K_2\rangle + (\epsilon - \delta)|K_1\rangle]/\sqrt{1+|\epsilon-\delta|^2}$$

where

$$|K_1\rangle = [|K^0\rangle + |\bar{K}^0\rangle]/\sqrt{2}$$

$$|K_2\rangle = [|K^0\rangle - |\bar{K}^0\rangle]/\sqrt{2}$$

and

$$|\bar{K}^0\rangle = CP|K^0\rangle.$$

The parameter δ specifies the *CPT*-violating part.

Estimates of δ are given below assuming the validity of the $\Delta S=\Delta Q$ rule.
See also THOMSON 95 for a test of *CPT*-symmetry conservation in K^0 decays using the Bell-Steinberger relation.

REAL PART OF δ

A nonzero value violates *CPT* invariance.

VALUE (units 10^{-4})	EVTS	DOCUMENT ID	TECN	COMMENT
2.9± 2.6±0.6	1.3M	3 ANGELOPO... 98F	CPLR	
• • • We do not use the following data for averages, fits, limits, etc. • • •				
2.4± 2.8		4 APOSTOLA... 99B	RVUE	
180 ± 200	6481	5 DEMIDOV 95		$K_{\ell 3}$ reanalysis

³ If $\Delta S=\Delta Q$ is not assumed, ANGELOPOULOS 98F finds $Re\delta=(3.0 \pm 3.3 \pm 0.6) \times 10^{-4}$.

⁴ APOSTOLAKIS 99B assumes only unitarity and combines CPLEAR and other results.

⁵ DEMIDOV 95 reanalyzes data from HART 73 and NIEBERGALL 74.

IMAGINARY PART OF δ

A nonzero value violates *CPT* invariance.

VALUE (units 10^{-5})	EVTS	DOCUMENT ID	TECN	COMMENT
- 0.2± 2.0		6 LAI	05A NA48	
• • • We do not use the following data for averages, fits, limits, etc. • • •				
2.4± 5.0		7 APOSTOLA... 99B	RVUE	
- 90 ± 290 ± 100	1.3M	8 ANGELOPO... 98F	CPLR	
2100 ± 3700	6481	9 DEMIDOV 95		$K_{\ell 3}$ reanalysis

⁶ LAI 05A values are obtained through unitarity (Bell-Steinberger relations), improving determination of η_{000} and combining other data from PDG and APOSTOLAKIS 99B.

⁷ APOSTOLAKIS 99B assumes only unitarity and combines CPLEAR and other results.

⁸ If $\Delta S=\Delta Q$ is not assumed, ANGELOPOULOS 98F finds $Im\delta=(-15 \pm 23 \pm 3) \times 10^{-3}$.

⁹ DEMIDOV 95 reanalyzes data from HART 73 and NIEBERGALL 74.

Re(y)

A non-zero value would violate *CPT* invariance in $\Delta S = \Delta Q$ amplitude. Re(y) is the following combination of K_{e3} decay amplitudes:

$$\text{Re}(y) = \text{Re} \left(\frac{A(\bar{K}^0 \rightarrow e^- \pi^+ \bar{\nu}_e)^* - A(K^0 \rightarrow e^+ \pi^- \nu_e)}{A(\bar{K}^0 \rightarrow e^- \pi^+ \bar{\nu}_e)^* + A(K^0 \rightarrow e^+ \pi^- \nu_e)} \right)$$

VALUE (units 10^{-3})	EVTS	DOCUMENT ID	TECN	
0.4 ± 2.5	13k	10	AMBROSINO 06E	KLOE
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.3 ± 3.1		11	APOSTOLA...	99B CPLR
10 They use the PDG 04 (web update) for the K_L^0 semileptonic charge asymmetry and PDG 04 (<i>CP</i> review, <i>CPT</i> NOT ASSUMED) for $\text{Re}(\epsilon)$.				
11 Constrained by Bell-Steinberger (or unitarity) relation.				

Re(x₋)

A non-zero value would violate *CPT* invariance in decay amplitudes with $\Delta S \neq \Delta Q$. x_- , used here to define Re(x₋), and x_+ , used below in the $\Delta S = \Delta Q$ section are the following combinations of K_{e3} decay amplitudes:

$$x_{\pm} = \frac{1}{2} \left(\frac{A(\bar{K}^0 \rightarrow \pi^- e^+ \nu_e)}{A(K^0 \rightarrow \pi^- e^+ \nu_e)} \pm \frac{A(K^0 \rightarrow \pi^+ e^- \bar{\nu}_e)^*}{A(\bar{K}^0 \rightarrow \pi^+ e^- \bar{\nu}_e)^*} \right).$$

VALUE (units 10^{-3})	EVTS	DOCUMENT ID	TECN	COMMENT
-0.8 ± 2.5	13k	12	AMBROSINO 06E	KLOE Tagged K_S^0
• • • We do not use the following data for averages, fits, limits, etc. • • •				
-0.5 ± 3.0		13	APOSTOLA...	99B CPLR Strangeness tagged
2 ± 13 ± 3	650k		ANGELOPO...	98F CPLR Strangeness tagged
12 Uses PDG 04 (web update) for the K_L^0 semileptonic charge asymmetry and $\text{Re}(\delta)$ from CPLEAR, ANGELOPOULOS 98F.				
13 Constrained by Bell-Steinberger (or unitarity) relation.				

$$|m_{K^0} - m_{\bar{K}^0}| / m_{\text{average}}$$

A test of *CPT* invariance. “Our Evaluation” is described in the “Tests of Conservation Laws” section. It assumes *CPT* invariance in the decay and neglects some contributions from decay channels other than $\pi\pi$.

VALUE	CL%	DOCUMENT ID	TECN	
< 10^{-18} (CL = 90%) OUR EVALUATION				
• • • We do not use the following data for averages, fits, limits, etc. • • •				
(- 3 ± 4) × 10^{-18}		14	ANGELOPO...	99B RVUE
14 ANGELOPOULOS 99B assumes only unitarity and combines CPLEAR and other results.				

$$(\Gamma_{K^0} - \Gamma_{\bar{K}^0}) / m_{\text{average}}$$

A test of *CPT* invariance.

VALUE	DOCUMENT ID	TECN	
(7.8 ± 8.4) × 10^{-18}	15	ANGELOPO...	99B RVUE
15 ANGELOPOULOS 99B assumes only unitarity and combines CPLEAR with other results. Correlated with $(m_{K^0} - m_{\bar{K}^0}) / m_{\text{average}}$ with a correlation coefficient of -0.95.			

TESTS OF $\Delta S = \Delta Q$ RULE

Re(x_+)

A non-zero value would violate the $\Delta S = \Delta Q$ rule in *CPT* conserving transitions. x_+ is defined above in the Re(x_-) section.

VALUE (units 10^{-3})	EVTS	DOCUMENT ID	TECN
-0.8±3.1 OUR AVERAGE			
-0.5±3.6	13k	16 AMBROSINO	06E KLOE
-1.8±6.1		17 ANGELOPO...	98D CPLR

16 $\text{Re}(x_+)$ can be shown to be equal to the following combination of rates:

$$\text{Re}(x_+) = \frac{1}{2} \frac{\Gamma(K_S^0 \rightarrow \pi e \nu) - \Gamma(K_L^0 \rightarrow \pi e \nu)}{\Gamma(K_S^0 \rightarrow \pi e \nu) + \Gamma(K_L^0 \rightarrow \pi e \nu)}$$

which is valid up to first order in terms violating *CPT* and/or the $\Delta S = \Delta Q$ rule.

17 Obtained neglecting *CPT* violating amplitudes.

K^0 REFERENCES

ABOUZAID	06	PRL 96 101801	E. Abouzaid <i>et al.</i>	(KTEV Collab.)
AMBROSINO	06E	PL B636 173	F. Ambrosino <i>et al.</i>	(KLOE Collab.)
LAI	05A	PL B610 165	A. Lai <i>et al.</i>	(CERN NA48 Collab.)
PDG	04	PL B592 1	S. Eidelman <i>et al.</i>	
LAI	03C	EPJ C30 33	A. Lai <i>et al.</i>	(CERN NA48 Collab.)
LAI	02	PL B533 196	A. Lai <i>et al.</i>	(CERN NA48 Collab.)
ANGELOPO...	01B	EPJ C22 55	A. Angelopoulos <i>et al.</i>	(CPLEAR Collab.)
ANGELOPO...	99B	PL B471 332	A. Angelopoulos <i>et al.</i>	(CPLEAR Collab.)
APOSTOLA...	99B	PL B456 297	A. Apostolakis <i>et al.</i>	(CPLEAR Collab.)
ANGELOPO...	98D	PL B444 38	A. Angelopoulos <i>et al.</i>	(CPLEAR Collab.)
Also		EPJ C22 55	A. Angelopoulos <i>et al.</i>	(CPLEAR Collab.)
ANGELOPO...	98E	PL B444 43	A. Angelopoulos <i>et al.</i>	(CPLEAR Collab.)
ANGELOPO...	98F	PL B444 52	A. Angelopoulos <i>et al.</i>	(CPLEAR Collab.)
Also		EPJ C22 55	A. Angelopoulos <i>et al.</i>	(CPLEAR Collab.)
DEMIDOV	95	PAN 58 968 From YAF 58 1041.	V. Demidov, K. Gusev, E. Shabalin	(ITEP)
THOMSON	95	PR D51 1412	G.B. Thomson, Y. Zou	(RUTG)
BARKOV	87B	SJNP 46 630 Translated from YAF 46	L.M. Barkov <i>et al.</i> 1088.	(NOVO)
BARKOV	85B	JETPL 42 138 Translated from ZETFP 42 113.	L.M. Barkov <i>et al.</i>	(NOVO)
BLATNIK	79	LNC 24 39	S. Blatnik, J. Stahov, C.B. Lang	(TUZL, GRAZ)
MOLZON	78	PRL 41 1213	W.R. Molzon <i>et al.</i>	(EFI+)
NIEBERGALL	74	PL 49B 103	F. Niebergall <i>et al.</i>	(CERN, ORSAY, VIEN)
HART	73	NP B66 317	J.C. Hart <i>et al.</i>	(CAVE, RHEL)
FOETH	69B	PL 30B 276	H. Foeth <i>et al.</i>	(AACH, CERN, TORI)
HILL	68B	PR 168 1534	D.G. Hill <i>et al.</i>	(BNL, CMU)
FITCH	67	PR 164 1711	V.L. Fitch <i>et al.</i>	(PRIN)
BALTAY	66	PR 142 932	C. Baltay <i>et al.</i>	(YALE, BNL)
BURNSTEIN	65	PR 138B 895	R.A. Burnstein, H.A. Rubin	(UMD)
KIM	65B	PR 140B 1334	J.K. Kim, L. Kirsch, D. Miller	(COLU)
CHRISTENS...	64	PRL 13 138	J.H. Christenson <i>et al.</i>	(PRIN)
CRAWFORD	59	PRL 2 112	F.S. Crawford <i>et al.</i>	(LRL)
ROSENFIELD	59	PRL 2 110	A.H. Rosenfeld, F.T. Solmitz, R.D. Tripp	(LRL)